**Problem 1** (5 points).

Consider a spherical cavity of radius  $a$  cut out of a conducting metal. A pure dipole  $p$  is placed at the center of the cavity. Find the potential inside the cavity.

**Solution**

Assume  $\vec{p} \parallel \hat{e}_3$ . Due to azimuthal symmetry we can expand the potential in Legendre polynomials

$$\Phi(r, \theta) = \sum_{l=0}^{\infty} (A_l r^l + B_l r^{-l-1})$$

The potential is a sum of the potential of pure dipole

$$\frac{p \cos \theta}{4\pi\epsilon_0 r^2}$$

and the potential of the induced charges on the conductor. At small  $r$  the potential due to the induced charges is  $\sim \frac{Q_{\text{ind}}}{4\pi\epsilon_0 a} \sim \text{const}$  so the singular behavior of the potential as  $r \rightarrow 0$  is coincides with the dipole potential. The corresponding coefficients  $B_l$  are thus

$$B_1 = \frac{p}{4\pi\epsilon_0}, \quad B_{l \neq 1} = 0$$

and the expansion in Legendre polynomials takes the form

$$\Phi(r, \theta) = \left( A_1 r + \frac{p}{4\pi\epsilon_0 r^2} \right) \cos \theta + \sum_{l \neq 1}^{\infty} A_l r^l P_l(\cos \theta)$$

From the boundary condition  $\Phi(a, \theta) = 0$  and orthogonality of Legendre polynomials it is clear that  $A_{l \neq 1} = 0$  and

$$A_1 a + \frac{p}{4\pi\epsilon_0 a^2} = 0 \quad \Rightarrow \quad A_1 = -\frac{p}{4\pi\epsilon_0 a^3}$$

The result for the potential is

$$\Phi(r, \theta) = \frac{p}{4\pi\epsilon_0} \left( -\frac{r}{a^3} + \frac{1}{r^2} \right) \cos \theta$$

**Problem 2** (6 points).

Find the Dirichlet Green function of Laplace equation for the interior of infinite cylinder with radius  $a$ .

**Solution**

Up to Eqs. (3.37) and (3.38) from “Chapter 3” file everything is the same as for infinite space. The difference is in the boundary condition for  $y_2(x')$ . For infinite space, we had  $y_2(x') \rightarrow 0$  as  $x' \rightarrow \infty$  so the proper choice was  $y_2(x') = K_m(x')$ . Now, the boundary condition is  $y_2(ka) = 0$  so we should take

$$y_2(x') = K_m(x') - \frac{K_m(ka)}{I_m(ka)} I_m(x')$$

The Wronskian  $W(y_1(x'), y_2(x')) = -\frac{1}{x'}$  is the same as for infinite space case since  $W(I(x'), I(x')) = 0$  so the Green function can be obtained from Eq. (3.7.45) by replacement of  $K_m(ks)$  by

$$L_m(ks) = K(ks) - \frac{K_m(ka)}{I_m(ka)} I_m(ks)$$

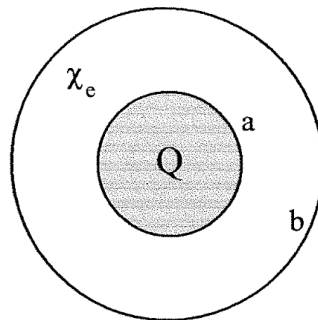
Finally, the Green function reads

$$G(\vec{r}, \vec{r}') = \frac{1}{\pi} \sum_{m=-\infty}^{\infty} \int_{-\infty}^{\infty} dk e^{ik(z-z')} I_m(|k|s_<) L_m(|k|s_>)$$

A quick check at  $a \rightarrow \infty$ : we get Eq. (3.7.45) since the additional term in  $L$  vanishes due to  $\frac{K_m(ka)}{I_m(ka)} \xrightarrow{a \rightarrow \infty} 0$ .

**Problem 1** (5 points).

A conducting sphere of radius  $a$  carries charge  $Q$ . It is surrounded by dielectric material of susceptibility  $\chi_e$ , out to radius  $b$ . Find the energy of this configuration.

**Solution**

Due to spherical symmetry, the electric field and displacement are radial. Due to Gauss' law, the electric displacement at  $r > a$  is

$$D(r) = \frac{Q}{4\pi r^2}$$

so the electric field is given by

$$E(r) = \frac{Q}{4\pi r^2 \epsilon} \theta(b-r) \theta(r-a) + \frac{Q}{4\pi r^2 \epsilon_0} \theta(r-b)$$

The energy of this setup is

$$W = \frac{1}{2} \int \vec{D} \cdot \vec{E} d^3x = 4\pi \frac{\epsilon}{2} \int_a^b dr r^2 E^2(r) + 4\pi \frac{\epsilon_0}{2} \int_b^\infty dr r^2 E^2(r) = \frac{Q}{8\pi \epsilon_0} \left[ \frac{1}{\chi_e} \left( \frac{1}{a} - \frac{1}{b} \right) + \frac{1}{b} \right]$$