

## 804 Midterm (16 points). 02/14/20

**Problem 1 .**

An iron sphere of radius  $R$  carries a charge  $Q$  and has a uniform magnetization  $\vec{M} = M\hat{e}_3$ .

It is initially at rest. Find

(a) Angular momentum stored in the fields

(b) If the sphere is demagnetized by heating (keeping  $\vec{M}$  uniform), by use of Faraday's law find the induced electric field, then find the torque induced by  $\vec{E}$  on the sphere, and finally the angular momentum imparted to sphere as  $M$  goes to zero.

Hint: The magnetic field outside the sphere is equal to the magnetic field of a pure dipole with  $\vec{m} = \frac{4}{3}\pi R^3 \vec{M}$

**Solution**

(a)

The electric and magnetic fields (at  $r > R$ ) are:

$$\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}, \quad \vec{B} = \frac{\mu_0}{4\pi r^3} (3\hat{r}(\vec{m} \cdot \hat{r}) - \vec{m}) \quad \Rightarrow \quad \vec{E} \times \vec{B} = \frac{Q\mu_0 m}{16\pi^2 \epsilon_0 r^5} \hat{e}_3 \times \hat{r}$$

where  $\vec{m} = \frac{4}{3}\pi R^3 \vec{M}$ .

The angular momentum stored in the fields is

$$\int_R^\infty dr r^2 \int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi \epsilon_0 \vec{r} \times (\vec{E} \times \vec{B}) = \int_R^\infty dr r^2 \int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi \frac{\mu_0 m Q}{16\pi^2 r^4} (\hat{e}_3 - (\hat{e}_3 \cdot \hat{r})\hat{r})$$

By symmetry, the angular momentum  $\vec{L}$  is collinear to  $\hat{e}_3$  so

$$\vec{L} = \hat{e}_3 \int_R^\infty dr r^2 \int_0^\pi d\theta \int_0^{2\pi} d\phi \frac{\mu_0 m Q}{16\pi^2 r^4} \sin^3\theta = \frac{\mu_0 Q}{6\pi R} \vec{m} = \frac{2}{9} \mu_0 Q R^2 M$$

(b)

If  $\vec{m} = m(t)\hat{e}_3$  the induced electric field takes the form

$$\vec{E} = -\frac{\partial \vec{A}}{\partial t} = -\frac{\partial}{\partial t} \frac{\mu_0 m(t) \hat{e}_3 \times \hat{r}}{4\pi r^2} = -\frac{\mu_0}{4\pi r^2} \hat{e}_3 \times \hat{r} \frac{dm}{dt} = -\frac{\mu_0}{4\pi r^2} \hat{\phi} \sin\theta \frac{dm}{dt}$$

The torque for the charge  $dq$  on the surface is

$$d\vec{\tau} = dq \vec{r} \times \vec{E} = dq \frac{\mu_0}{4\pi R} \hat{\theta} \sin\theta \frac{dm}{dt}, \quad \hat{r} \times \hat{\phi} = -\hat{\theta}$$

Again, from symmetry we know that  $\vec{\tau} \parallel \hat{e}_3$  so

$$d\tau_3 = -dq \frac{\mu_0}{4\pi R} \sin^2\theta \frac{dm}{dt}$$

Since the charge is distributed uniformly over the surface of the sphere  $dq = \frac{Q}{4\pi R^2} R^2 \sin\theta d\theta d\phi$  and the total torque is

$$\tau_3 = - \int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi \sin^2\theta \frac{\mu_0 Q}{16\pi^2 R} \frac{dm}{dt} = - \frac{\mu_0 Q}{6\pi R} \frac{dm}{dt}$$

The angular momentum imparted to sphere is

$$\vec{m} = \hat{e}_3 \int_0^\infty \tau_3(t) dt = - \frac{\mu_0 Q}{6\pi R} \hat{e}_3 \int_0^\infty \frac{dm}{dt} dt = \frac{\mu_0 Q}{6\pi R} m \hat{e}_3 = \frac{2}{9} \mu_0 Q R^2 M \hat{e}_3$$

**Problem 2** (6 points)

A circular wire of radius  $a$  lies in  $x, y$  plane with the center at the origin. It carries current

$$I(t) = I_0 \theta(t) = \begin{cases} 0 & t < 0 \\ I_0 & t \geq 0 \end{cases}$$

Find the magnetic field on the  $z$  axis (at  $x, y = 0$ ).

**Solution:**

Generalization of Bio-Savart law for time-dependent currents has the form (formula (6.56) from *Jackson* or (6.3.23) from Chapter 6 of lecture notes)

$$\vec{B}(\vec{x}, t) = \frac{\mu_0}{4\pi} \int d^3x' \left[ \frac{\vec{J}(\vec{x}', t_{\text{ret}})}{R^2} + \frac{\dot{\vec{J}}(\vec{x}', t_{\text{ret}})}{cR} \right] \times \hat{e}_R$$

where  $t_{\text{ret}} = t - \frac{R}{c}$  and  $R = |\vec{x} - \vec{x}'|$ . Rewriting it for linear currents, one gets

$$\vec{B}(\vec{x}, t) = \frac{\mu_0}{4\pi} \int \left[ \frac{I(\vec{x}', t_{\text{ret}})}{R^2} + \frac{\dot{I}(\vec{x}', t_{\text{ret}})}{cR} \right] d\vec{l}' \times \hat{e}_R$$

By symmetry, the magnetic field has only  $z$  component. Moreover, for a point on  $z$ -axis the retarded time is the same for all points on the wire and is equal to  $t - \frac{R}{c} = t - \frac{\sqrt{z^2 + a^2}}{c}$ . Also,  $I(t) = I_0 \theta(t)$  and  $\dot{I}(t) = I_0 \delta(t)$  so we get

$$\begin{aligned} B_3(\vec{x}, t) &= \frac{\mu_0}{4\pi} \int \left[ \frac{I_0}{z^2 + a^2} \theta\left(t - \frac{\sqrt{z^2 + a^2}}{c}\right) + \frac{I_0}{c\sqrt{z^2 + a^2}} \delta\left(t - \frac{\sqrt{z^2 + a^2}}{c}\right) \right] (d\vec{l}' \times \hat{e}_R)_3 \\ &= \frac{\mu_0 I_0}{4\pi(z^2 + a^2)} \theta\left(t - \frac{\sqrt{z^2 + a^2}}{c}\right) \int (d\vec{l}' \times \hat{e}_R)_3 + \frac{\mu_0 I_0}{4\pi c \sqrt{z^2 + a^2}} \delta\left(t - \frac{\sqrt{z^2 + a^2}}{c}\right) \int (d\vec{l}' \times \hat{e}_R)_3 \end{aligned}$$

From geometry

$$(d\vec{l}' \times \hat{e}_R)_3 = dl' \frac{a}{\sqrt{a^2 + z^2}} \Rightarrow \int (d\vec{l}' \times \hat{e}_R)_3 = \frac{2\pi a^2}{\sqrt{z^2 + a^2}}$$

and therefore

$$B(z, t) = \frac{\mu_0 I_0 a^2 \hat{e}_3}{2(z^2 + a^2)^{3/2}} \theta\left(t - \frac{\sqrt{z^2 + a^2}}{c}\right) + \frac{\mu_0 I_0 a^2 \hat{e}_3}{2c(z^2 + a^2)} \delta\left(t - \frac{\sqrt{z^2 + a^2}}{c}\right)$$

Note that the first term is equal to the result for constant current  $I_0$  multiplied by  $\theta\left(t - \frac{\sqrt{z^2 + a^2}}{c}\right)$  (see e.g. Eq. (5.38) from Griffiths textbook). The second part is a burst of radiation coming from turning on the current.

### Problem 3.

Find the reflection coefficient for the circularly polarized electromagnetic wave incident on a plane between two linear media at Brewster's angle. (For simplicity, take  $\mu' = \mu$ ).

### Solution:

The circularly polarized wave have the form

$$\vec{E} = \frac{1}{\sqrt{2}}(\hat{e}_1 + i\hat{e}_2)E_0^i e^{ik(z \cos \theta_i + x \sin \theta_i)}$$

(we've assumed that  $\vec{k}$  lies in the  $XZ$  plane).

Let us consider the components  $E_0^i \hat{e}_1$  (electric field in the plane of incidence) and  $iE_0^i \hat{e}_2$  (electric field normal to the plane of incidence) separately.

At  $\theta_i = \theta_B$  we have no reflected wave for the component  $E_0^i \hat{e}_1$  (in the plane of incidence). For  $iE_0^i \hat{e}_2$  (normal to the plane of incidence) we have

$$E_0^r = \left(\frac{i}{\sqrt{2}}E_0^i\right) \frac{\sin(\theta_T - \theta_i)}{\sin(\theta_T + \theta_i)}$$

Since at  $\theta_i = \theta_B$  we have  $\sin \theta_B = \cos \theta_T = n'/\sqrt{n^2 + n'^2}$  and  $\cos \theta_B = \sin \theta_T = n/\sqrt{n^2 + n'^2}$ , we obtain

$$E_0^r = \frac{i}{\sqrt{2}}E_0^i \frac{\sin(\theta_T - \theta_B)}{\sin(\theta_T + \theta_B)} = \frac{i}{\sqrt{2}}E_0^i \frac{n^2 - n'^2}{n^2 + n'^2}$$

Therefore, the intensity of the reflected wave is proportional to  $\frac{1}{2}\left(E_0^i \frac{n^2 - n'^2}{n^2 + n'^2}\right)^2$  and the reflection coefficient is

$$R = \frac{(n^2 - n'^2)^2}{2(n^2 + n'^2)^2}$$