Phys. 807 — Statistical Mechanics

HW8

Problem 1.

In the neighborhood of z = 1 the following expansion may be obtained

$$g_{5/2}(z) = 2.363\nu^{3/2} + 1.342 - 2.612\nu - 0.730\nu^2 + \dots$$
(1)

where $\nu = -\ln z$. From this, the corresponding expansions for $g_{3/2}, g_{1/2}$ and $g_{-1/2}$ may be obtained by the recursion formula $g_{n-1} = -\partial g_n / \partial \nu$. Using this expansion show that for the ideal Bose gas the discontinuity of $\partial C_V / \partial T$ at $T = T_c$ is given by

$$\left(\frac{\partial}{\partial T} \frac{C_V}{Nk}\right)_{T \to T_c^-} - \left(\frac{\partial}{\partial T} \frac{C_V}{Nk}\right)_{T \to T_c^+} = \frac{c}{T_c} \ .$$

and find constant c (numerically).

Solution

We start from Eq. (10.57) from the lecture notes rewritten as

$$\frac{C_V(T)}{Nk_B} = \begin{cases} \frac{15}{4} \frac{v}{\lambda^3} g_{5/2}(\alpha^*) - \frac{9}{4} \frac{g_{3/2}(\alpha^*)}{g_{1/2}(\alpha^*)} & T > T_c \\ \frac{15}{4} \frac{v}{\lambda^3} g_{5/2}(1) & T < T_c \end{cases}$$
(2)

where we used Eq. (10.34): $\frac{v}{\lambda^3} = g_{3/2}(\alpha^*)$. Now

$$\left(\frac{\partial}{\partial T}\frac{C_V}{Nk}\right)_{T\to T_c^-} = \left.\frac{15}{4}g_{5/2}(1)v\frac{\partial}{\partial T}\frac{1}{\lambda^3}\right|_{T=T_c}$$

and

$$\left(\frac{\partial}{\partial T} \frac{C_V}{Nk} \right)_{T \to T_c^+}$$

$$= \left. \frac{15}{4} g_{5/2}(1) v \frac{\partial}{\partial T} \frac{1}{\lambda^3} \right|_{T=T_c} + \left. \frac{\partial \alpha^*}{\partial T} \left(\frac{15}{4\alpha^*} - \frac{9}{4\alpha^*} + \frac{9}{4\alpha^*} \frac{g_{3/2}(\alpha_*)g_{-1/2}(\alpha_*)}{g_{1/2}^2(\alpha_*)} \right) \right|_{T \to T_c}$$

where we used $\alpha \frac{\partial}{\partial \alpha} g_n(\alpha) = g_{n-1}(\alpha)$. Next, $\frac{d\alpha}{dT}$ is known from Eq. (10.55) from the lecture notes

$$\frac{d\alpha}{dT} = -\frac{3}{2T} g_{3/2}(\alpha) \frac{\alpha}{g_{1/2}(\alpha)}$$
(3)

so we get

$$\left(\frac{\partial}{\partial T} \frac{C_V}{Nk} \right)_{T \to T_c^-} - \left(\frac{\partial}{\partial T} \frac{C_V}{Nk} \right)_{T \to T_c^+}$$

$$= - \frac{9}{4T} \frac{g_{3/2}(\alpha^*)}{g_{1/2}(\alpha^*)} \Big|_{T \to T_c} - \frac{27}{8T} \frac{g_{3/2}^2(\alpha^*)g_{-1/2}(\alpha^*)}{g_{1/2}^3(\alpha^*)} \Big|_{T \to T_c}$$

Since $g_{3/2}(1) = 2.612$ and $g_{-1/2}(1) = \infty$ the first term vanishes as $T \to T_c$. To estimate the second term, we need the expansion of $g_{-1/2}$ and $g_{1/2}$ near $\alpha^* = 1 \ (\equiv \nu \to 0)$ which can be obtained from Eq. (1):

$$g_{1/2}(\alpha^*) = \frac{3}{4}(2.363)\nu^{-1/2} - 2(0.730) + \dots$$

$$g_{-1/2}(\alpha^*) = \frac{3}{8}(2.363)\nu^{-3/2} + \dots$$

We see that $\frac{g_{-1/2}(\alpha^*)}{g_{1/2}^3(\alpha^*)}\Big|_{\alpha^*\to 1} = \frac{8}{9(2.363)^2}$ so

$$\left(\frac{\partial}{\partial T} \frac{C_V}{Nk}\right)_{T \to T_c^-} - \left(\frac{\partial}{\partial T} \frac{C_V}{Nk}\right)_{T \to T_c^+} = -\frac{3}{T_c} \left(\frac{g_{3/2}(1)}{2.363}\right)^2$$

and therefore

$$c = -3(\frac{2.612}{2.363})^2 \simeq -3.665$$

Problem 2.

Show that the equation of state of the ideal Bose gas phase can be written in the form of a virial expansion, i.e.,

$$\frac{Pv}{kT} = 1 + c_1 \left(\frac{\lambda^3}{v}\right) + c_2 \left(\frac{\lambda^3}{v}\right)^2 - \dots$$

and find constants c_1 and c_2 .

Solution

The equation of state of an ideal Bose gas is given by Eq. (10.39) from the lecture notes (at $T > T_c$)

$$\frac{Pv}{k_BT} = \frac{v}{\lambda^3} g_{5/2}(\alpha_*) = \frac{v}{\lambda^3} [\alpha^* + \frac{\alpha_*^2}{2^{5/2}} + \frac{\alpha_*^3}{3^{5/2}} + \dots]$$

On the other hand

$$\frac{\lambda^3}{v} = g_{3/2}(\alpha_*) = \alpha^* + \frac{\alpha_*^2}{2^{3/2}} + \frac{\alpha_*^3}{3^{3/2}} + \dots$$
(4)

 \mathbf{SO}

$$\frac{Pv}{k_BT} = \frac{\alpha^* + \frac{\alpha^2_*}{2^{5/2}} + \frac{\alpha^3_*}{3^{5/2}} + \dots}{\alpha^* + \frac{\alpha^2_*}{2^{3/2}} + \frac{\alpha^3_*}{3^{3/2}} + \dots} = 1 - \frac{\alpha_*}{4\sqrt{2}} + \alpha^2_* (\frac{1}{16} - \frac{2}{9\sqrt{3}}) + O(\alpha^3_*)$$
(5)

In order to get the equation of state in the form of virial expansion we must invert Eq. (4)

$$\alpha_* = \frac{\lambda^3}{v} \left[1 - \frac{1}{2\sqrt{2}} \frac{\lambda^3}{v} + O(\frac{\lambda^6}{v^2})\right]$$

and substitute it into Eq. (5). We get

$$\frac{Pv}{k_BT} = 1 - \frac{1}{4\sqrt{2}}\frac{\lambda^3}{v} + \frac{\lambda^6}{v^2}\left[\frac{1}{8} - \frac{2}{9\sqrt{3}}\right] + O(\frac{\lambda^9}{v^3})$$
$$c_1 = -\frac{1}{4\sqrt{2}} \quad \text{and} \quad c_2 = \frac{1}{8} - \frac{2}{9\sqrt{3}}$$

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