

**Problem 1.**

A surface charge glued over the (non-conducting) sphere of radius  $R$  produces the potential

$$\phi(x, y, z) = \frac{V_0}{2R^2}(2z^2 - x^2 - y^2)$$

inside the sphere. Find the potential outside the sphere and the surface charge density.

**Solution**

Expand the potential in Legendre polynomials. The potential inside is

$$\phi^{\text{in}}(r, \theta) = \frac{V_0 r^2}{2R^2}(3 \cos^2 \theta - 1)$$

The potential is continuous so outside of the sphere we have

$$\phi^{\text{out}}(r, \theta) = \frac{V_0 R^3}{2r^3}(3 \cos^2 \theta - 1)$$

The surface charge density is

$$\sigma(r, \theta) = \epsilon_0(E^{\text{out}} - E^{\text{in}})|_{r=R} = \epsilon_0 \left( \frac{\partial \phi^{\text{in}}}{\partial r} - \frac{\partial \phi^{\text{out}}}{\partial r} \right) \Big|_{r=R} = \frac{5\epsilon_0 V_0}{2R}(3 \cos^2 \theta - 1)$$

**Problem 2.**

Find the Dirichlet Green function of Laplace equation for the interior of infinite cylinder with radius  $a$ .

**Solution**

Up to Eqs. (3.37) and (3.38) from "Chapter 3" file everything is the same as for infinite space. The difference is in the boundary condition for  $y_2(x')$ . For infinite space, we had  $y_2(x') \rightarrow 0$  as  $x' \rightarrow \infty$  so the proper choice was  $y_2(x') = K_m(x')$ . Now, the boundary condition is  $y_2(ka) = 0$  so we should take

$$y_2(x') = K_m(x') - \frac{K_m(ka)}{I_m(ka)} I_m(x')$$

The Wronskian  $W(y_1(x'), y_2(x')) = -\frac{1}{x'}$  is the same as for infinite space case since  $W(I(x'), I(x')) = 0$  so the Green function can be obtained from Eq. (3.7.45) by replacement of  $K_m(ks)$  by

$$L_m(ks) = K(ks) - \frac{K_m(ka)}{I_m(ka)} I_m(ks)$$

Finally, the Green function reads

$$G(\vec{r}, \vec{r}') = \frac{1}{\pi} \sum_{m=-\infty}^{\infty} \int_{-\infty}^{\infty} dk e^{ik(z-z')} I_m(|k|s_<) L_m(|k|s_>)$$

A quick check at  $a \rightarrow \infty$ : we get Eq. (3.7.45) since the additional term in  $L$  vanishes due to  $\frac{K_m(ka)}{I_m(ka)} \xrightarrow{a \rightarrow \infty} 0$ .

**Problem 3.**

A half-space below  $z = 0$  is filled with dielectric with susceptibility  $\chi_e > 1$ . A pure dipole  $\vec{p} = p\hat{e}_1$  is located at distance  $d$  above the dielectric. Find the force acting on the dipole. Is it attractive or repulsive?

**Solution**

The image dipole is  $\vec{p}' = -p \frac{\chi_e - 1}{\chi_e + 1}$ . A general formula for the force acting on the dipole  $\vec{p}$  located at  $\vec{x}$  due to the dipole  $\vec{p}'$  located at  $\vec{x}'$  is

$$\vec{F}(\vec{x}) = \frac{1}{4\pi\epsilon_0} \nabla_x \left\{ \frac{3(\vec{p}' \cdot (\vec{x} - \vec{x}'))(\vec{p} \cdot (\vec{x} - \vec{x}')) - (\vec{p} \cdot \vec{p}')|\vec{x} - \vec{x}'|^2}{|\vec{x} - \vec{x}'|^5} \right\}$$

In our case this reduces to

$$\vec{F}(\vec{x}) = -\frac{1}{4\pi\epsilon_0}(\vec{p} \cdot \vec{p}')\nabla_x \frac{1}{|\vec{x} - \vec{x}'|^3} = -\frac{3p^2}{4\pi\epsilon_0} \frac{\chi_e - 1}{\chi_e + 1} \frac{\vec{x} - \vec{x}'}{|\vec{x} - \vec{x}'|^5} = -\frac{3p^2}{64d^4\pi\epsilon_0} \frac{\chi_e - 1}{\chi_e + 1} \hat{e}_3$$

Thus, the force is attractive, the direction is down, and the magnitude is  $\frac{3p^2}{4\pi\epsilon_0} \frac{\chi_e - 1}{\chi_e + 1} (2d)^{-4}$

All problems have equal weight.